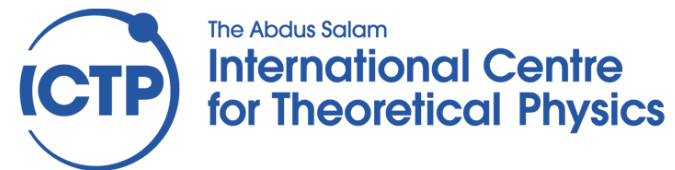


# Energy conservation and Jarzynski equality are incompatible for quantum work

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In collaboration with



Alberto Imparato



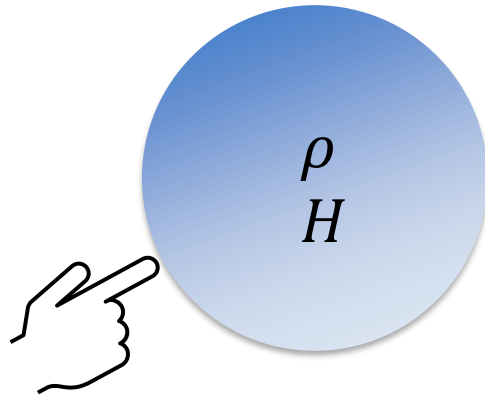
# Plan of the talk

- Minimal requirements on work
- Linear measurement schemes
- Nonlinear measurement schemes
- The ultimate No-Go
- Summary

**arXiv:2104.09364**

# Paradigm

Closed ( $\Leftrightarrow$  thermally isolated) system is driven by external coherent fields



# Main problem

- Can the fluctuations of work be observed?
- If we observe *some* statistics, how do we know if it is work?

# Our approach

- Take the most general representation of work statistics quantum mechanics can offer
- Establish a minimal set of physical requirements work statistics must satisfy
- Check whether these requirements are consistent with the laws of quantum mechanics

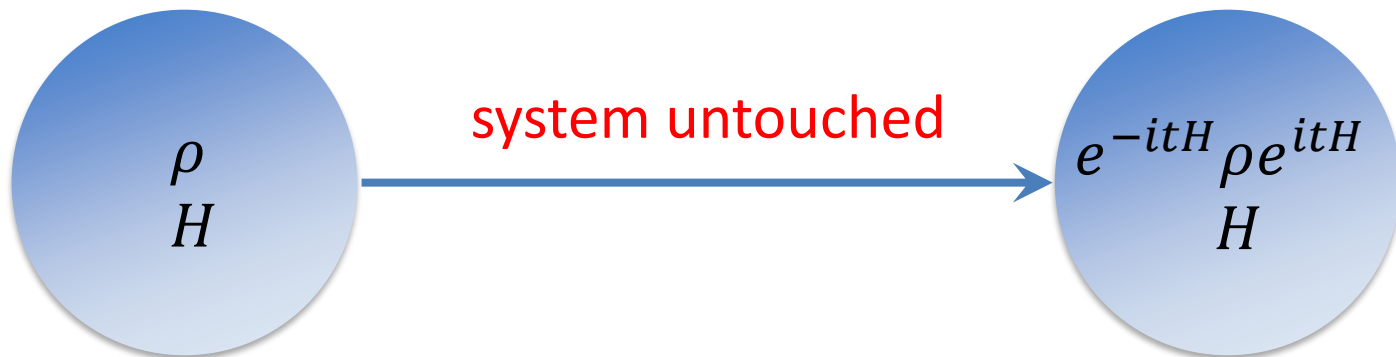
Inspired by the approach taken in

*Perarnau-Llobet, Bäumer, Hovhannisyan, Huber, & Acin, PRL **118**, 070601 (2017)*

➤ Minimal requirements on work

# What is to be expected from work

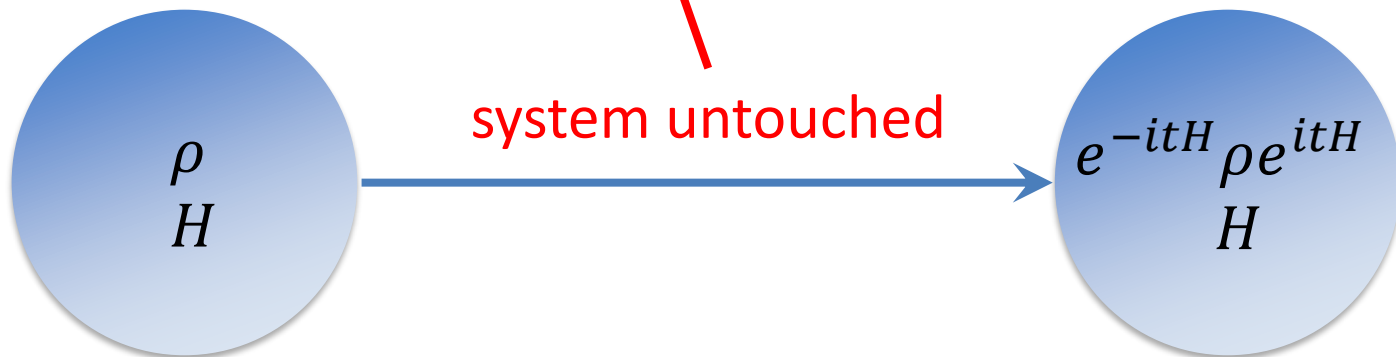
(A) Nothing happens





# What is to be expected from work

(A) Nothing happens  $\Rightarrow$  deterministically zero work is exchanged



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$$p(W) = \delta(W) \quad \forall \rho$$

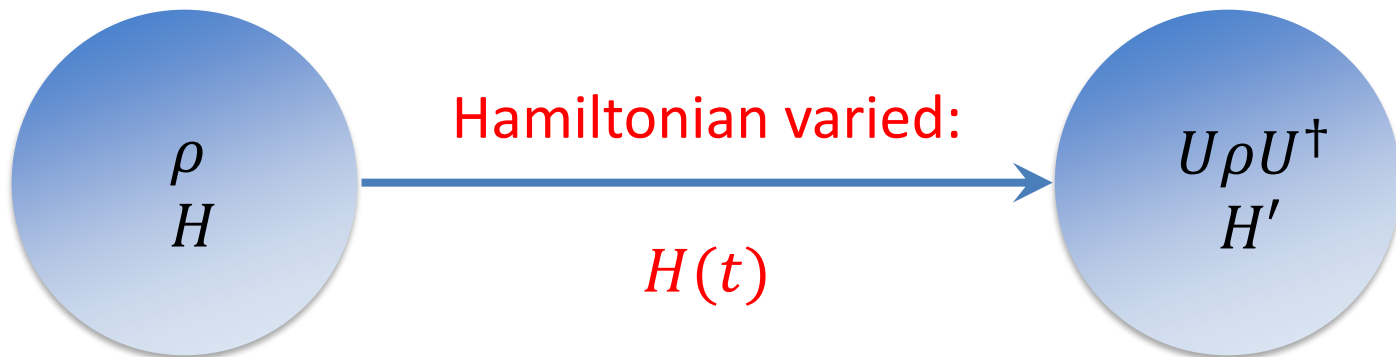


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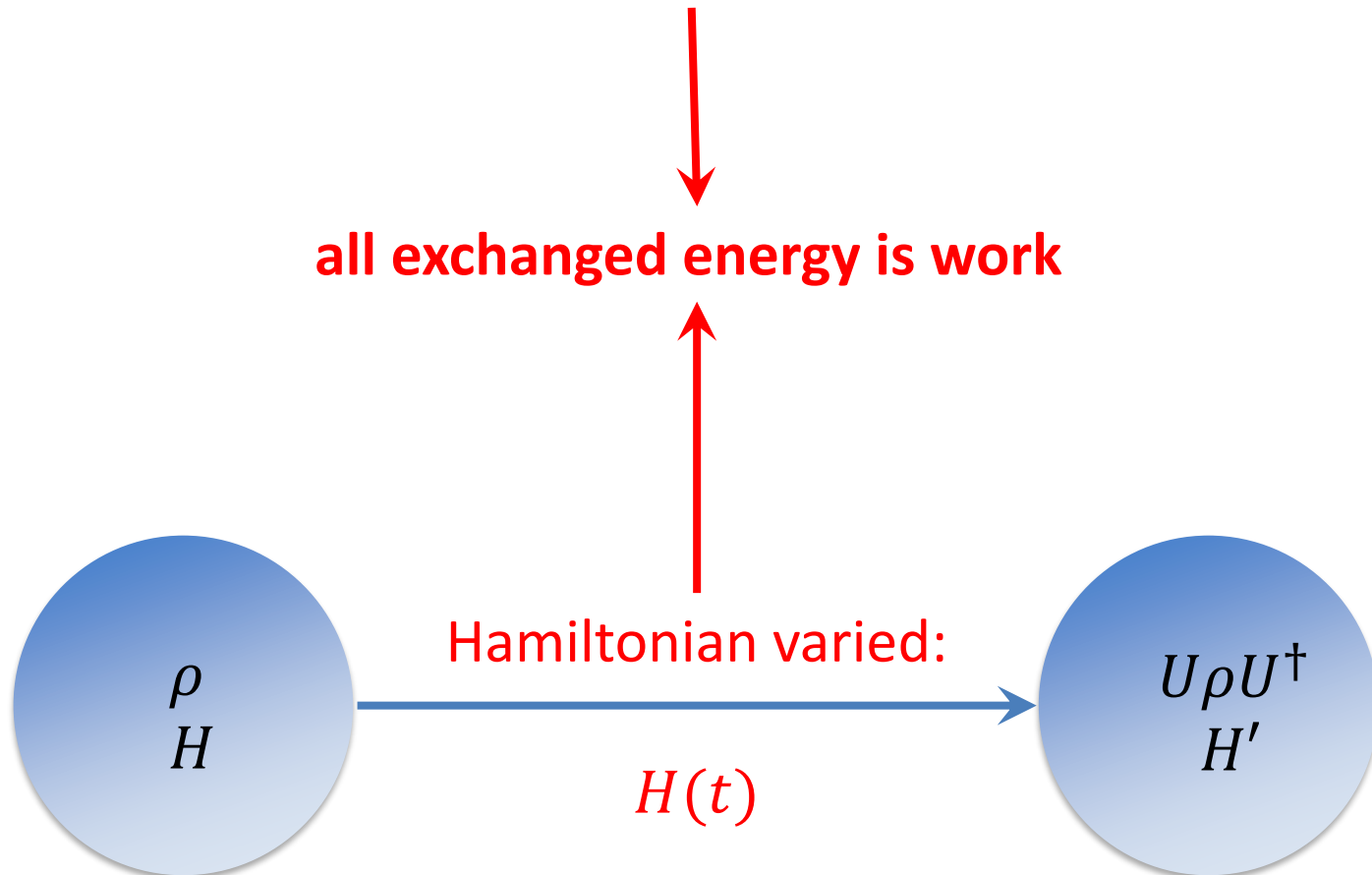


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**both points ensure that the observed statistics respect the energy conservation**



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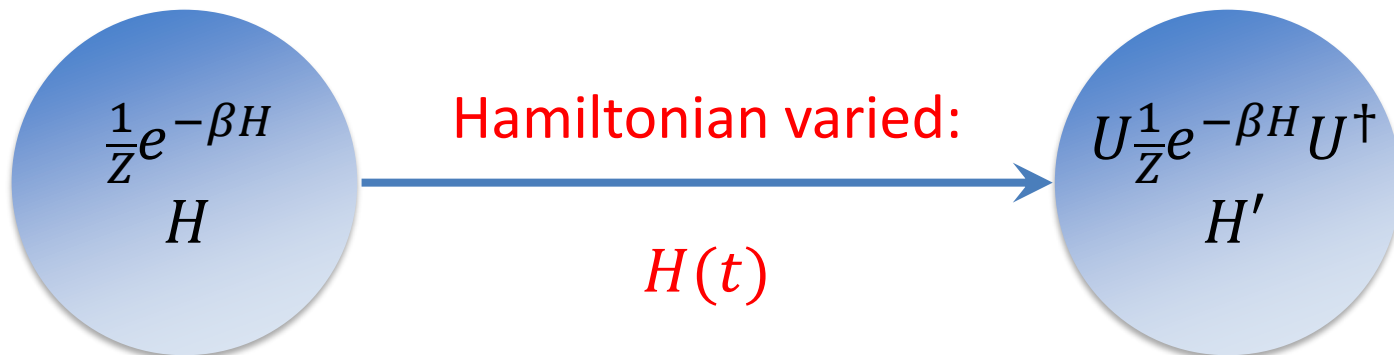
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$$\langle e^{-\beta W} \rangle = e^{-\beta \Delta F} \quad \forall \beta$$



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**no a priori need for this, but it is a nice-to-have bridge to the second law and to the classical realm**





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**plus most major work-measuring protocols satisfy it**



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# Existing schemes and $(A)$ , $(A')$ , $(B)$

## Two-projective-measurement (TPM) scheme

*Kurchan*, arXiv:cond-mat/0007360

*Tasaki*, arXiv:cond-mat/0009244

## Operator of work

*Allahverdyan & Nieuwenhuizen*, PRE **71**, 066102 (2005)

## Bohmian mechanics approach

*Sampaio, Suomela, Ala-Nissila, Anders, & Philbin*, PRA **97**, 012131 (2018)

## Bayesian network approach

*Micadei, Landi, & Lutz*, arXiv:2103.14570 [quant-ph]

## “Single measurement” approaches

*Deffner, Paz, & Zurek*, PRE **94**, 010103(R) (2016)

*Gherardini, Belenchia, Paternostro, & Trombettoni*, arXiv:2106.06461 [quant-ph]

## “Resource-theoretical” approach

*Åberg*, PRX **8**, 011019 (2018)

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**for coherent processes, these papers do not construct a random variable for work, so they do not address the problem here**

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## “Resource-theoretical” approach

**while quasiprobability is a useful notion, constructing a quasiprobability distribution for work does not address the problem of measurability**

## Quasiprobabilistic approach

*Allahverdyan*, Phys. Rev. E 90, 032137 (2014)

*Solinas & Gasparinetti*, Phys. Rev. E 92, 042150 (2015)

*Miller & Anders*, New J. Phys. 19 062001 (2017)

The examples show that all the pairs of  $(A)$ ,  $(A')$ , and  $(B)$  are jointly satisfiable.

What about jointly satisfying all three?

➤ Linear schemes

# Linearity



$$p_W(\lambda_1\rho_1 + \lambda_2\rho_2) = \lambda_1 p_W(\rho_1) + \lambda_2 p_W(\rho_2)$$

$$\forall \lambda_1, \lambda_2 \geq 0, \quad \lambda_1 + \lambda_2 = 1$$

# Linearity



$$p_W(\lambda_1\rho_1 + \lambda_2\rho_2) = \lambda_1 p_W(\rho_1) + \lambda_2 p_W(\rho_2)$$



$$p_W(\rho) = \text{Tr}(\rho M_W)$$

constitute a POVM:

$$M_W \geq 0 \ \& \ \sum_W M_W = I$$

# Linearity



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$$p_W(\rho) = \text{Tr}(\rho M_W)$$

$M_W$  are state-independent!

## Result 1

If a linear scheme satisfies (A') and (B) for all thermal states:

$$\sum_W W \operatorname{Tr}(\tau_\beta M_W) = \operatorname{Tr}(\tau_\beta U^\dagger H' U) - \operatorname{Tr}(\tau_\beta H)$$



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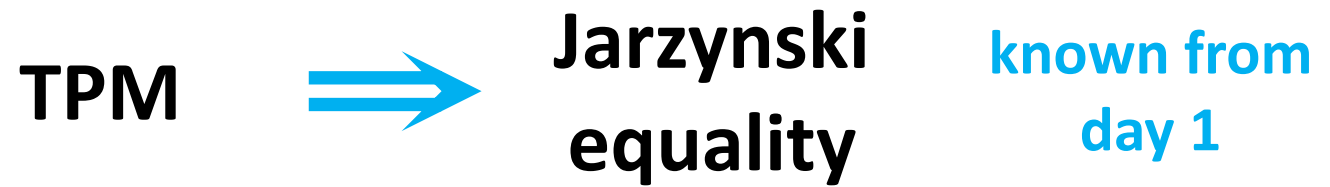
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then it coincides with the TPM scheme:

$$M_W = M_W^{\text{TPM}}$$



**our  
Result 1**

**TPM**



**Jarzyński  
equality**

**known from  
day 1**

**our  
Result 1**

**TPM**



**Jarzynski  
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**known from  
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Thus, linear schemes that satisfy  $(B)$  for thermal states cannot satisfy  $(A')$  for all states.

➤ Nonlinear schemes

# Nonlinear schemes

Any conceivable distribution for work can be thought of as one generated by a POVM performed on the system:

$$p_W = \text{Tr}(\rho M_W)$$

In general,  $p_W$  will **not** be linear in  $\rho$ :

$$M_W = M_W(\rho)$$

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In general,  $p_W$  will **not** be linear in  $\rho$ :

$$M_W = M_W(\rho)$$

Can this extra freedom allow reconciling  $(A)$ ,  $(A')$ , and  $(B)$ ?



# Artificial nonlinear protocol

## Artificial Protocol

- When  $[\rho, H] = 0$ , perform TPM
- When  $[\rho, H] \neq 0$ , perform the operator of work scheme

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Conditions  $(A)$ ,  $(A')$ , and  $(B)$  are satisfied simultaneously!

# Artificial nonlinear protocol

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- When  $[\rho, H] = 0$ , perform TPM
- When  $[\rho, H] \neq 0$ , perform the operator of work scheme



Conditions  $(A)$ ,  $(A')$ , and  $(B)$  are satisfied simultaneously!

**However,**  
the corresponding measurement is not continuous wrt  $\rho$

# The ultimate No-Go

## Result 2

No measurement exists for which all  $M_W(\rho)$  are continuous in  $\rho$  and the conditions (A), (A'), and (B) are satisfied simultaneously.

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No measurement exists for which all  $M_W(\rho)$  are continuous in  $\rho$  and the conditions (A), (A'), and (B) are satisfied simultaneously.

## Result 3

If none of the outcomes  $\{W\}$  goes to 0 as  $U \rightarrow \mathbb{I}$ , then the scheme cannot simultaneously satisfy (A) and (A') — the two expressions of energy conservation.

# Take-home messages

- For a linear measurement, satisfying Jarzynski's equality is equivalent to being the TPM scheme.
- Continuous nonlinear measurements cannot reconcile energy conservation and Jarzynski's equality.
- If a scheme has work outcomes that do not depend on the unitary evolution operator, then it cannot respect energy conservation.

**arXiv:2104.09364**